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Synthesis, X-ray and optical characterizations of two new oxysulfides: LaInS₂O and La₅In₃S₉O₃^{$\frac{A}{A}$}

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Abstract

Two new compounds, LaInS₂O and La₅In₃S₉O₃ were synthesized in the La–In–S–O quaternary system. Both compounds crystallize in the orthorhombic system with lattice constants $a = 20.5421(6)$ \AA , $b = 14.8490(4)$ \AA , $c = 3.9829(1)$ \AA for LaInS₂O, and $a = 4.1018(1)$ Å, $b = 26.833(1)$ Å, $c = 16.023(1)$ Å for La₅In₃S₉O₃. The structure of La₅In₃S₉O₃ was solved from single-crystal Xray data, in the space group Pbcm, with $Z = 4$; it is built from three-atom-thick (100)_{NaCl} layers interleaved with fluorite-type ribbons, and is closely related to the structures of the known lanthanum and indium compounds $La_{10}In_{6}S_{17}O_{6}$ and $La_{4}In_{5}S_{13}$. Both compounds LaInS₂O and La₅In₃S₉O₃ exhibit a yellow color; measurement of their optical gaps gave 2.73 and 2.60 eV, respectively. \odot 2003 Elsevier Inc. All rights reserved.

Keywords: Oxysulfide; Indium; Lanthanum; Crystal structure; Optical properties

1. Introduction

For a few years the increased importance of transparent conducting compounds for application technologies such as low emissivity windows or thin film photovoltaics led to a renewed interest for this kind of materials [\[1\]](#page-5-0). A considerable effort was done to improve existing materials as well as for the search of new materials. Owing to the limitations of existing materials the need for new p-type transparent conducting compounds is especially high. In association with ntype materials, the discovery of efficient p-type materials would open the way to transparent electronics.

Recently, a few new p-type transparent conducting materials have been reported such as $SrCuO₂$ [\[1\]](#page-5-0) and LaCuSO [\[2\]](#page-5-0). LaCuSO presents a composite type structure that results from an alternated stacking of fluorite type layers $[La_2O_2]$ with anti-fluorite type layers $[Cu_2S_2]$ [\[3\].](#page-5-0) For this compound, band structure calculations show that the valence band is made of Cu and S states while the bottom of the conduction band is made

of La states [\[4\].](#page-5-0) The p-type conductivity is achieved by replacing part of La by Sr. According to the composite character of the structure, this doping does not alter the conducting $\left[\text{Cu}_2\text{S}_2\right]$ layer, which could favor a high electrical conductivity.

We have started the search for new p-type transparent conducting compounds with a composite type structure. As In³⁺ is a d^{10} cation with the d energy levels not too far from those of Cu, we have focused our attention on the La–In–S–O system. This paper reports on the chemical, X-ray diffraction and optical characterizations of two new compounds $LaInS_2O$ and $La₅In₃S₉O₃$ that we found during our exploration. The crystal structure of $La₅In₃S₉O₃$ has been solved and will be discussed in relation with the homeotypic structure of $La_{10}In_{6}S_{17}O_{6}$ [\[5\]](#page-5-0).

2. Experimental

2.1. Synthesis and chemical analysis

LaInS₂O and La₅In₃S₉O₃ were obtained during an attempt to synthesize an hypothetical composite compound $La_2In_2S_4O_2$ built from the stacking of both [$La₂O₂$] and [InS₂] slabs. Powders of $La₂S₃$ (preliminary

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obtained by sulfurization of $La₂O₃$ under a gas flow of H_2S at 1200°C), La₂O₃, and In₂S₃ were weighed in a 1:2:3 ratio, respectively. The mixture was loaded into a silica tube sealed under vacuum ($p = 2 \times 10^{-2}$ atm) and heated at a rate of 100° C h⁻¹ up to 800°C for 12 h. This reaction yielded a pure yellow powder. The powder was ground and subsequently divided in two equal amounts. Each of them was sealed in an evacuated silica tube and reheated at 850° C or 950° C with a small amount of iodine $(<5 \text{ mg cm}^{-3})$ to favor crystallization. In both cases, a mixture of a yellow powder, black octahedral crystals, and yellow needle-like crystals was obtained.

Chemical analysis by means of energy dispersive Xray spectrometry (EDXS) revealed that the black octahedral crystals have the composition In_2S_3 . For the two yellow phases the chemical analysis confirmed the presence of oxygen and we found the following atomic percentages for the other elements $(at\%)$:

- yellow powder La 24.7% , In 23.1% , S 52.2% , in good agreement with the initial composition $LalnS₂O;$
- yellow needle-like crystals La 30.3% , In 14.9% , S 54.8%, in good agreement with the formulation $La₅In₃S₉O₃$ deduced from the single-crystal structure determination (at%: La 29.4, In 17.7, S 52.9).

2.2. Optical measurements

Optical measurements were performed with a Leitz spectrophotometer (MPV-SP), coupled with a metallographic microscope, working in the visible and close IR range (400–800 nm). For $La_5In_3S_9O_3$, two spectra in polarized transmitted light were recorded on a single crystal, with polarization directions perpendicular and parallel to the elongation of the crystal $(a\text{-axis})$. On the other hand, the diffuse reflectance spectra of $LaInS₂O$ and $In₂S₃$ were recorded from powders.

2.3. X-ray powder diffraction

The X-ray powder diffraction pattern of $LaInS₂O$ was recorded on a Bragg–Brentano Siemens D5000 diffractometer. Prior to measurement the yellow powder was ground, passed through a $20 \mu m$ mesh, and placed on a flat sample holder.

2.4. Data collection and structure refinement

A single crystal of $La_5In_3S_9O_3$ with a parallelepipedic shape (dimensions $0.005 \times 0.04 \times 0.1$ mm³) was mounted on a Nonius Kappa CCD diffractometer, using MoKa radiation ($\lambda = 0.71073 \text{ Å}$). The crystal-to-detector distance was 40.0 mm. 559 frames were collected with a rotation angle per frame of 0.8° and an exposure time of 92 s. The diffraction pattern was consistent with the

orthorhombic system. We could refine the unit-cell parameter values: $a = 4.1018(1)$ Å, $b = 26.833(1)$ Å, and $c = 16.023(1)$ Å, from 27,948 reflections with $2.91^{\circ} < \theta < 34.97^{\circ}$.

Intensities were corrected for absorption $(\mu = 16.73 \text{ mm}^{-1})$ using the face indexed option. The structure was solved with the use of the direct methods (SHELXS), and refinements with subsequent difference Fourier were done with SHELXL (SHELXTL software package [\[6\]\)](#page-5-0). The structure solution was easily obtained. We had to split an In position into two sites In(2) and In(3) that are too close to be occupied simultaneously. We have then constrained the sum of the occupancies of In(2) and In(3) to be equal to 1, and their thermal parameters to be identical. In the final stage of refinement, i.e., with all atoms (except oxygen atoms) refined anisotropically, the reliability factors converged to $R_1 = 0.0447$ and $wR_2 = 0.0811$ for 1627 observed reflections $(I>2\sigma(I))$, restricted to $2\theta_{\text{max}} = 60^{\circ}$ $(\sin \theta / \lambda \approx 0.7)$, and 96 parameters. The highest and deepest peaks $(4.71 e \text{ Å}^{-3})$; $-2.83 e \text{ Å}^{-3}$) on the final Fourier difference map were closed to $S(2)$ (1.43 Å) and La(3) (1.39 Å), respectively. For comparison, the height of the highest residual peak corresponds to about 1/6 of that of an O atom. Details on data collection and structure refinements are given in Table 1. Atomic coordinates and equivalent isotropic displacement parameters are listed in [Table 2](#page-2-0); coefficients of the anisotropic displacement parameters are listed in

Table 1

Crystal data and structure refinement for $La₅In₃S₉O₃$

Empirical formula Formula weight $(g \text{ mol}^{-1})$ Temperature Wavelength Crystal system Space group	$La_5In_3S_9O_3$ 1375.5 293(2) K 0.71073 Å Orthorhombic Phcm
Unit-cell dimensions	$a = 4.1018(1)$ Å
	$b = 26.8325(5)$ Å $c = 16.0232(3)$ Å
Volume	$1763.54(6)$ Å ³
Z	4
Density (calculated)	5.191 mg/m ³
Absorption coefficient	16.734 mm^{-1}
F(000)	2405
Crystal size	$0.005 \times 0.04 \times 0.1$ mm ³
θ range for data collection	$2.54 - 30.00^{\circ}$
Index ranges	$-5 \le h \le 5$, $-37 \le k \le 37$, $-22 \le l \le 22$
Reflections collected	17520
Independent reflections	2660 $[R_{\text{int}} = 0.0903]$
Completeness to $\theta = 30.00^{\circ}$	99.8%
Refinement method	Full-matrix least-squares on F^2
Data/restraints/parameters	2660/0/96
Goodness-of-fit on F^2	1.145
<i>R</i> indices $[1627 I > 2\sigma(I)]$	$R_1 = 0.0447$, $wR_2 = 0.0811$
R indices (all data)	$R_1 = 0.0951$, $wR_2 = 0.0916$
Extinction coefficient	0.00004(1)
Largest diff. peak and hole	4.72 and $-2.83 \text{ e} \text{ Å}^{-3}$

Table 2 Atomic coordinates and equivalent isotropic displacement parameters $(\AA^2 \times 10^3)$ for La₅In₃O₃S₉

Atom	Site	S.o.f	\boldsymbol{x}	υ	Z	$U_{\rm eq}$
La(1)	8e		0.2505(2)	0.5210(1)	0.6198(1)	5(1)
La(2)	4d		$-0.2498(4)$	0.6134(1)	3/4	5(1)
La(3)	8e		$-0.2530(3)$	0.6256(1)	0.4990(1)	6(1)
In(1)	4d		$-0.2600(5)$	0.3820(1)	3/4	11(1)
In(2)	8e	0.852(2)	$-0.4284(3)$	0.2449(1)	0.6288(1)	11(1)
In(3)	8e	0.148(2)	$-0.070(2)$	0.2446(2)	0.6287(5)	$11(-)$
S(1)	8e		$-0.2403(13)$	0.3316(1)	0.6226(2)	10(1)
S(2)	4d		0.242(2)	0.4348(1)	3/4	11(1)
S(3)	8e		0.2475(13)	0.5534(1)	0.4465(2)	7(1)
S(4)	4c		$-0.778(2)$	1/4	1/2	12(1)
S(5)	8e		$-0.7549(13)$	0.6582(1)	0.6202(2)	8(1)
S(6)	4d		0.214(2)	0.2342(2)	3/4	18(1)
O(1)	4d		0.252(5)	0.5606(4)	3/4	7(2)
O(2)	8e		$-0.248(3)$	0.5673(3)	0.6155(5)	7(1)

 U_{eq} is defined as one-third of the trace of the orthogonalized U^{ij} tensor.

Table 3 Anisotropic displacement parameters ($\AA^2 \times 10^3$) for La₅In₃S₉O₃

	U^{11}	I^{122}	U^{33}	I^{123}	I^{13}	U^{12}
La(1)	4(1)	6(1)	6(1)	$-1(1)$	0(1)	0(1)
La(2)	4(1)	6(1)	5(1)	0	$\overline{0}$	0(1)
La(3)	4(1)	8(1)	7(1)	1(1)	$-1(1)$	0(1)
In(1)	9(1)	11(1)	12(1)	0	θ	$-1(1)$
In(2)	16(1)	7(1)	10(1)	1(1)	1(1)	0(1)
In(3)	16(1)	7(1)	10(1)	1(1)	1(1)	0(1)
S(1)	12(1)	6(1)	11(1)	0(1)	5(3)	0(2)
S(2)	11(2)	10(2)	12(2)	0	$\mathbf{0}$	0(3)
S(3)	7(1)	7(1)	8(1)	$-1(1)$	1(2)	$-1(2)$
S(4)	8(3)	17(2)	12(2)	$-2(2)$	$\mathbf{0}$	θ
S(5)	8(1)	8(1)	8(1)	0(1)	1(3)	$-3(2)$
S(6)	25(4)	18(2)	9(2)	0	θ	2(3)

The anisotropic displacement factor exponent takes the form: $-2\pi^2[h^2a^{*2}U^{11} + \cdots + 2hka^*b^*U^{12}].$

Table 3. Selected interatomic distances are given in [Table 4](#page-3-0).

Nota: Considering the tendency for the X-coordinates to be near 1/4 or 3/4, a refinement of the structure in the space group *Cmcm* was performed; it yielded a satisfying reliability factor but too many observed weak reflections violating the C centering mode $(h + k = 2n)$ had to be rejected.

2.5. Results and discussion

In the Ln–Ga–S–O system, numerous compounds were reported such as $LaGaS_2O$ [\[7\],](#page-5-0) La_3GaS_3O [\[8\]](#page-6-0), $(LaO)_4Ga_{1,33}S_4$ and $(LnO)_4Ga_2S_5$ [\[9\];](#page-6-0) the two latter have layered structures formed by alternation of a $[Ln_2O_2]$ slab and a $[Ga_xS_y]$ slab. On the contrary, before our investigation, the only known compound in the quaternary system La–In–S–O was $La_{10}In_{6}S_{17}O_{6}$. This compound was synthesized from a mixture of La_2O_2S and In_2S_3 (1:1 ratio) at 1500 K by Gastaldi et al. [\[5\]](#page-5-0). At lower temperature $(<1223 \text{ K})$ our investigation of the La–In–S–O quaternary system did not yield $La_{10}In_6$ $S_{17}O_6$ but the two title compounds, namely LaInS₂O and $La₅In₃S₉O₃$.

[Fig. 1](#page-4-0) gives the X-ray powder pattern of $LaInS₂O$. We could index the pattern in the orthorhombic crystal system (see [Table 5](#page-4-0)), and we have refined the cell parameters as $a = 20.5421(6)$ Å, $b = 14.8490(4)$ Å, $c = 3.9829(1)$ A using the profile matching mode as implemented in the FullProf software [\[10\].](#page-6-0) This compound seems not to be isotypic with $LaGaS₂O$, and we were unable to solve the structure neither using direct methods nor by using a trial and error method.

On the other hand, we were able to solve the structure of $La_5In_3S_9O_3$ which was obtained as single crystals. [Fig. 2](#page-4-0) shows the structure of this compound projected onto the (b, c) plane. We observe a segregation of the sulfur and oxygen atoms. Only sulfur atoms coordinate indium atoms, while both oxygen and sulfur atoms coordinate lanthanum atoms. In(1) and $In(2)/In(3)$ atoms are found in tetrahedral coordination of sulfur atoms with In–S distances ranging from 2.28 to 2.52 Å (see [Table 4](#page-3-0)). The sum of the occupancies for In(2) and In(3) was constrained to be equal to 1, which forbids, statistically, the simultaneous occupation of both sites, and then the too short $In(2)$ – $In(3)$ contacts of 1.472 and 2.631 Å, along the **a** direction; the filling is 85.2% and 14.8% for $In(2)$ and $In(3)$, respectively.

As In(2) and In(3) atoms could be ordered along the a-direction, we have checked for the existence of a superstructure along this direction. Precession photographs obtained on the Kappa CCD diffractometer did not reveal any superstructure along a^* that would indicate a long range order for the distribution of the filled $In(2)$ and $In(3)$ sites.

La(1) is seven-fold coordinated $(1 O(1), 2 O(2), 3 S(3),$ and $1 S(2)$, with four atoms in an equatorial plane $(2 O(2), 2 S(3))$ and three other atoms $(O(1), S(2), S(3))$

Table 4 Bond lengths (A) and bond valence sums for $La₅In₃S₉O₃$

Atoml	Atom2	Sym.	Distance	Bond val.
La(1)	O(1)		2.340(5)	0.64
La(1)	O(2)		2.392(11)	0.55
La(1)	O(2)	#1	2.405(11)	0.53
La(1)	S(3)		2.910(3)	0.48
La(1)	S(3)	#2	3.048(4)	0.33
La(1)	S(3)	#3	3.059(4)	0.32
La(1)	S(2)		3.115(3)	0.28
La(1)	La(2)	#1	3.8336(14)	$\overline{}$
La(1)	La(3)	#1	3.9710(11)	
				$\Sigma = 3.13$
La(2)	O(2)		2.485(8)	0.43
La(2)	O(2)	#4	2.485(8)	0.43
La(2)	O(1)	#5	2.489(18)	0.42
La(2)	O(1)		2.497(17)	0.42
La(2)	S(5)	#1	3.145(4)	0.26
La(2)	S(5)	#6	3.145(4)	0.25
La(2)	S(5)	#4	3.172(4)	0.24
La(2)	S(5)		3.172(4)	0.24
La(2)	S(6)	#7	3.244(5)	0.20
				$\Sigma = 2.89$
La(3)	O(2)		2.436(7)	0.49
La(3)	S(3)	#5	2.943(4)	0.44
La(3)	S(3)		2.946(4)	0.44
La(3)	S(5)	#1	2.951(4)	0.43
La(3)	S(5)		2.962(4)	0.42
La(3)	S(1)	#2	3.034(4)	0.34
La(3)	S(1)	#9	3.071(4)	0.31
La(3)	S(4)	#9	3.3399(7)	0.15
La(3)	La(1)	#5	3.9712(11)	
				$\Sigma = 3.02$
In(1)	S(1)		2.451(3)	0.78
In(1)	S(1)	#4	2.451(3)	0.78
In(1)	S(2)	#5	2.486(7)	0.71
In(1)	S(2)		2.499(7)	0.69
				$\Sigma = 2.96$
In(2)	In(3)		1.472(7)	$\overline{}$
In(2)	S(5)	#10	2.448(3)	0.79
In(2)	S(6)	#5	2.451(5)	0.78
In(2)	S(1)		2.454(3)	0.78
In(2)	S(4)		2.518(4)	0.65
In(2)	In(3)	#5	2.630(7)	÷,
				$\Sigma=3.00$
In(3)	S(6)		2.283(9)	1.23
In(3)	S(4)	#1	2.388(8)	0.93
In(3)	S(5)	#10	2.431(7)	0.83
In(3)	S(1)		2.440(7)	0.81
In(3)	In(2)	#1	2.630(7)	
				$\Sigma = 3.80$

Symmetry transformations used to generate equivalent atoms: $\frac{\text{#1x+1}}{\text{#1x+1}}$ y, z ; #2-x,-y+1,-z+1; #3-x+1,-y+1,-z+1; #4x,y,-z+3/2; #5x-1,y,z; $\frac{\text{#6x+1}}{y}$, -z+3/2; $\frac{\text{#7-x}}{y}$ +1/2, -z+3/2; $\frac{\text{#9-x-1}}{y}$, -y+1, -z+1; $\frac{\text{#10-x-1}}{y}$ $y-1/2, z$.

in a perpendicular plane. This arrangement is usually described as a seven-octahedron with one apex split into two positions. La(2) is nine-fold coordinated $(2 \text{ O}(1), 2)$ $O(2)$, 4 $S(5)$, and 1 $S(6)$), which corresponds to a tricapped $(O(2), O(2), S(6))$ trigonal prism. Finally,

La(3) is eight-fold coordinated $(O(2), 2 S(1), 2 S(3), 2$ $S(5)$, and 1 $S(4)$), which is identified as the classical bicapped $(O(2), S(4))$ trigonal prism. The La–S and La–O distances range from 2.91 to 3.34 and 2.34 to 2.50 Å , respectively, as usually observed in lanthanum oxysulfides.

As shown in [Fig. 3a](#page-5-0), the crystal structure of $La₅In₃S₉O₃$ can be described as the combination of two types of building blocks. The first type is a distorted derivative of a three-atom-thick $(100)_{\text{NaCl}}$ layer, with composition $La_{1.5}In_{1.5}S₃$. The second building block is a two-atom-thick, two-tetrahedra-large ribbon of the fluorite type, with composition $[La_5O_3]$. The fluorite building blocks share La atoms on each side of the 2D infinite NaCl-type building block to form a zigzag pattern. Taking into account common La atoms, the connection of both building blocks corresponds to 2D infinite $La₅In₃S₆O₃$ layers. These layers alternate regularly along the b-direction with a corrugated S monoatomic layer.

This description of the structure with building blocks is very useful to highlight the similarity with the hightemperature homeotypic structure of $La_{10}In_6S_{17}O_6$. Indeed, this compound can be described in the same way, but here the layered NaCl-type building block acts as a mirror plane regarding to the fluorite building block, changing the space group from *Pbcm* to *Immm* (see [Fig. 3b\)](#page-5-0). Another difference between the title compound and $La_{10}In_{6}S_{17}O_{6}$ is the lack of one S atom in the inner atomic slab of the layered NaCl-type building block of $La_{10}In_6S_{17}O_6$. This leads to an indium pairing with an In–In distance of 2.86 A close to two times the covalent radius of indium (3.00 Å) . For $La₅In₃S₉O₃$, the corresponding In atoms (i.e. the two adjacent In(2)/In(3) couples along the c-direction) are separated by one S(6) atom (see [Fig. 2\)](#page-4-0), and the resulting In–In distance is greater than 3.8 Å . Finally both compounds $La_5In_3S_9O_3$ and $La_{10}In_6S_{17}O_6$ are derivatives of $La_4In_5S_{13}$ [\[11\]](#page-6-0). This last structure can be described as composed of a layer derived from a threeatom-thick $(100)_{NaCl}$ slab, with composition $La_4In_5S_9$ $[=(La₂InS₃)(In₃S₃)(La₂InS₃)],$ alternating with a sulfur mono-atomic plane (see [Fig. 3c\)](#page-5-0).

Bond valence sums calculated according to the method described in reference [\[12–13\]](#page-6-0) confirm the trivalent states of all La atoms, and In(1) as well as In(2) in $La_5In_3S_9O_3$ (see Table 4). However, the bond valence calculated for In(3) is much higher than 3. As the In(3) site is occupied only at 15% , the In(3)–S distances are shorter than expected, which explains this discrepancy. For the S-deficient derivative $La_{10}In_{6-}$ $S_{17}O_{16}$ [\[5\]](#page-5-0), we have calculated similar bond valence sums for La (3.02–3.23) and In (2.90–3.12) except for one indium atom, for which the valence sum is 1.70. This is in accordance with the formation of $(In₂)⁴⁺$ pairs in this compound.

Fig. 1. Observed X-ray diffraction pattern recorded for LaInS2O (open circle). In this picture the black line represents the calculated intensities obtained by profile matching, and the bottom curve shows the difference between experimental and calculated intensities.

Table 5 Positions $(2\theta^{\circ})$ and relative intensities of the main diffraction peaks measured for LaInS₂O

hkl	2θ	$I_{\rm exp}$	hkl	2θ	$I_{\rm exp}$
220	14.70	12.8	521	33.55	53.2
240	25.49	13.6	721	40.01	13.9
121	25.72	43.8	061	42.99	24.1
600	26.00	56.9	161	43.22	20.7
221	26.82	57.2	821	43.63	10.2
340	27.30	100.0	002	45.50	32.8
321	28.55	40.1	461	46.64	17.1
440	29.66	10.3	940	46.67	12.7
421	30.83	94.3	602	53.16	14.5
250	31.32	15.2	342	53.88	24.3
331	31.62	13.7	580	54.17	14.2
720	32.78	14.4			

[Fig. 4](#page-5-0) shows the absorbance spectra of $In₂S₃$, LaInS₂O and La₅In₃S₉O₃ obtained after a Kubelka– Munk transformation [\[14\]](#page-6-0). The crossing point between the base line along the energy axis and the extrapolated line of the absorption edge gives the optical gap. For In_2S_3 and $LaInS_2O$, the optical gaps are respectively 2.24 and 2.73 eV. For the compound $La_5In_3S_9O_3$, the values are 2.60 eV (perpendicular to *a*) and 2.65 eV (parallel to a). These gap values explain the yellow color of both compounds. As shown in [Fig. 5](#page-5-0), optical gaps increase with the ratio $O/(O + S)$ for In or La or mixed La/In compounds. According to the crystal structure study of $La_5In_3S_9O_3$, indium atoms are only bounded

Fig. 2. Projection of the structure of $\text{La}_3\text{In}_3\text{So}_3$ onto the (b, c) plane.

with S atoms in this compound while La is bonded to O and S atoms. We could expect the La oxychalcogenide part to give rise to an optical transition intermediate between those of La_2S_3 and La_2O_2S . This transition should be higher than the observed transitions for $La₅In₃S₉O₃$ and $LaInS₂O$. Consequently, the observed

Fig. 3. Modular description of the crystal structures of (a) $La_5In_3S_9O_3$, (b) $La_1_0In_6S_17O_6$ and (c) $La_4In_5S_13$. The fluorite-type ribbons and the threeatom-thick (100)_{NaCl} layers are indicated on this figure as well as the S-monoatomic layer. Small black circle=O; large black circle=S; grey circle=In; empty circle=La.

Fig. 4. Diffuse reflectance spectra of In_2S_3 (a) and $LaInS_2O$ (d) obtained after a Kubelka-Munk transformation at room temperature $(K/S$ vs. E). Polarized transmitted light spectra of La₅In₃S₉O₃ measured with the polarization directions perpendicular (b) and parallel (c) to the a -axis.

Fig. 5. Optical gaps of $La_5In_3S_9O_3$ and $La_{10}In_6S_{17}O_6$ compared to those of In_2S_3 , In_2O_3 , La_2S_3 and La_2O_2S . The gaps are plotted versus the ratio of oxygen over all anions $O/(O + S)$.

transitions may be related to the In/S part and the optical gap increase observed from In_2S_3 to $La_5In_3S_9O_3$ or $LaInS₂O$ could result from the inductive effect of the associated, more electropositive, La cation.

3. Conclusions

During our investigation of the La–In–S–O quaternary system we could not find a transparent compositelike oxysulfide. However, we have found two new yellow compounds, namely $LaInS₂O$ and $La₅In₃S₉O₃$. The structure of $La_5In_3S_9O_3$ exhibits $La_{1.5}In_{1.5}S_3$ infinite [NaCl] type layers and $[La₅O₃]$ fluorite-type ribbons. Both building blocks are also present in the structure of the lanthanum and indium oxysulfide $La_{10}In_{6}S_{17}O_{6}$. It is most likely that the structure determination of LaInS₂O would also reveal such blocks. We have measured the optical gaps for $LaInS_2O$ and $La₅In₃S₉O₃$ and we have found 2.73 and 2.60 eV, respectively. These gaps are intermediate between those of In_2S_3 and In_2O_3 . As the indium atoms are connected preferentially to sulfur atoms, it seems difficult to get a transparent compound in the La–In–S–O system.

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